Fast spin initialization of an electron trapped in a gated nanowire quantum dot

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Presentation plan



- 2 Spin initialization How is it done?
- 3 The proposed nanodevice
- 4 How we do it? The Rashba spin-orbit interaction

5 Brief summary

Quantum computer qubit

One of the many possible quantum computer implementations is based on the spin of an electron trapped in a semiconductor nananostructure.



(a) Elzerman et al., Science 333, 1269 (2011)



(b) Nowack, Nature 430, p. 431–435 (2004)

Advantages: easily integrable with classical electronics, scalability.



The concept of a qubit based on the spin of an electron trapped in semiconductor quantum dots first appeared in a (20 years old) paper by Loss and DiVincenzo:

Loss, DiVincenzo, Phys. Rev. A 57, 120 (1998)

Question: Why don't we have a working quantum processor? **Answer:** Due to the lack of a fast and effective spin initialization method.

A fast quantum interface between spin qubits of different codes, A. Noiri et al., arXiv:1804.04764 (2018).

Initialization in a strong magnetic field

The only currently available electron spin initialization method (in described systems) is based on splitting the spin levels in a strong magnetic field.





We apply the magnetic field and wait until the electron relaxes to the ground state. The method works but it is slow and requires very low temperatures $(10 \,\mathrm{mK})$.

Initialization in a strong magnetic field

In one of experiments done at the University of Delft (Holand) all necessary qubit operations were achieved.

K. C. Nowack et al., Science 333, 1261 (2011).



Spin initialization with fidelity of the order of 99% takes a couple of ns. The estimated spin coherence time is only about 20 ns. We need a faster method!

Proposed nanodevice

Let us create a nanodevice capable of trapping a single electron and performing all the necessary quantum operations on its spin. Setting the spin in a desired direction (initialization) **must be very fast**.



Bednarek, Pawłowski, Górski, Skowron, Phys. Rev. Applied 11, 034012 (2019)

Proposed nanodevice

- To the gates U_{1..7} we apply voltages forming an approximately parabolic confinement potential.
- A single electron is inserted into the wire.
- The electron relaxes to the ground state.
- The electron is interacted with by using the Rashba spin-orbit interaction (SOI) and voltages, applied to the gates.
- The spin-orbit interaction is a relativistic correction to the Hamiltonian but in most semiconductors it is strong enough to facilitate spin rotations.



In a bulk semiconductor with an electric field **E** present the SOI Hamiltonian is given by

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$$\hat{\mathbf{H}}_{so} = \frac{\alpha_{so}|\mathbf{e}|}{\hbar} \left(\mathbf{E} \times \hat{\mathbf{p}} \right) \cdot \hat{\sigma}. \qquad (1) \quad \left| \begin{array}{c} \mathbf{e}_{\mathbf{x}} \\ \mathbf{e}_{\mathbf{y}} \\ \mathbf{e}_$$

Our nanodevice constricts motion to the x axis, thus the Hamiltonian reduces to

$$\hat{H}_{\rm so} = \frac{\alpha_{\rm so}|\boldsymbol{e}|}{\hbar} \left(\boldsymbol{E}_{\boldsymbol{z}} \sigma_{\boldsymbol{y}} - \boldsymbol{E}_{\boldsymbol{y}} \sigma_{\boldsymbol{z}} \right) \hat{\boldsymbol{p}}_{\boldsymbol{x}}. \tag{2}$$

If we apply a strong electric field along the y axis, namely E_y , the second term dominates

$$\hat{H}_{\rm so} = -\frac{\alpha_{\rm so}|\boldsymbol{e}|}{\hbar} \boldsymbol{E}_{\boldsymbol{y}} \sigma_{\boldsymbol{z}} \hat{\boldsymbol{p}}_{\boldsymbol{x}}.$$
(3)

ton gate

The full Hamiltonian assumes the form

$$\hat{H} = \frac{\hat{p}_x^2}{2m^*} + U(x) + \hat{H}_{\rm so}.$$
 (4)

If we assume a **parabolic confinement potential** and **apply an electric field** E_y we obtain

$$\hat{H} = \frac{\hat{p}_x^2}{2m^*} + \frac{1}{2}m^*\omega^2 \hat{x}^2 - \frac{\alpha_{\rm so}|\boldsymbol{e}|}{\hbar} \boldsymbol{E}_y \sigma_z \hat{p}_x.$$
(5)

The eigenstates of this hamiltonian can be easily found in the momentum representation.

Eigenvalue equation:

$$\left[\frac{\hat{p}_{x}^{2}}{2m^{*}}+\frac{1}{2}m^{*}\omega^{2}\hat{x}^{2}-\frac{\alpha_{\rm so}|\boldsymbol{e}|}{\hbar}E_{y}\langle\sigma_{z}\rangle\hat{p}_{x}\right]\Psi=E\Psi.$$
(6)

The ground state:

$$\Psi = \sqrt{\frac{2\beta}{\pi}} e^{-\beta x^2} e^{iqx} \tag{7}$$

where:
$$\beta = \frac{m^* \omega}{2\hbar}$$
 and $q = \frac{m^* \alpha_{so} |e| E_y}{\hbar^2} \langle \sigma_z \rangle.$

The Gaussian wavepacket is multiplied by a plane wave, so **it should move**, **but it does not**¹ until we turn off the electric field.

¹wlewo.exe wprawo.exe

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We can also induce electron's movement if we proceed in the other way around. First, the electron is relaxed to the ground state and then the electric field E_y is turned on abruptly. The resulting electron motion is reversed².

If we use a spinor representation, the spin-degenerate ground state is given by:

$$\Psi_{\uparrow} = \sqrt{\frac{2\beta}{\pi}} \begin{pmatrix} 1\\0 \end{pmatrix} e^{-\beta x^2} e^{iqx} \qquad \Psi_{\downarrow} = \sqrt{\frac{2\beta}{\pi}} \begin{pmatrix} 0\\1 \end{pmatrix} e^{-\beta x^2} e^{-iqx}$$
(8)
where: $\beta = \frac{m^* \omega}{2\hbar}$ and $q = \frac{m^* \alpha_{\rm so} |e| E_y}{\hbar^2}.$

Given this representation we can construct a wavefunction of any possible spin orientation (e.g. equally-weighted linear combination).

²spinup.exe spindown.exe

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- No initial electric field $(E_y = 0)$.
- The initial state with spin oriented along the x axis given by

$$oldsymbol{\Psi}_{\mathsf{x}} = rac{1}{\sqrt{2}} \left(oldsymbol{\Psi}_{\uparrow} + oldsymbol{\Psi}_{\downarrow}
ight) = rac{1}{\sqrt{2}} egin{pmatrix} 1 \ 1 \end{pmatrix} e^{-eta x^2}.$$

• After the *E_y* field is turned on, the spin-up and spin-down components behave differently³.

(9)

³separacja.exe

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Second qubit

After the spin separation we obtain an interesting state. We can define the, so-called **charge qubit** which assigns qubit states to the presence of an electron in the left or in the right half of the nanodevice. Let us denote these states as:

$$|L\rangle = \Psi_L(x)$$
 $|R\rangle = \Psi_R(x).$ (10)

The presented state can be written down as

$$\Psi_{\mathsf{K}} = \begin{pmatrix} \Psi_{L}(x) \\ 0 \end{pmatrix} + \begin{pmatrix} 0 \\ \Psi_{R}(x) \end{pmatrix} = |\uparrow\rangle |L\rangle + |\downarrow\rangle |R\rangle.$$
(11)





 $\left|\uparrow\right\rangle\left|L\right\rangle+\left|\downarrow\right\rangle\left|R\right\rangle$

A similar state was obtained experimentally by D. Wineland in ion traps⁴.

The only difference: spin states replaced by two **internal (electronic)** states of the trapped ion.

Wineland obtained the so-called Schrödinger's cat states

$$\left|\uparrow\right\rangle \left|\odot\right\rangle + \left|\downarrow\right\rangle \left|\odot\right\rangle$$

J. Pawłowski, M. Górski, G. Skowron, and S. Bednarek Phys. Rev. B **96**, 115308 (2017)

⁴Monroe, Meekhof, King, Wineland, Science **272**, 1131 (1996)

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Entangled states

Our state is an example of **entangled states**⁵.

 $\left|\uparrow\right\rangle\left|L\right\rangle+\left|\downarrow\right\rangle\left|R\right\rangle$

A two-qubit state is called entangled if it is not possible to express it as a product of two one-qubit states.

Example 1:

$$\left|\uparrow\right\rangle \left|L
ight
angle + \left|\downarrow\right\rangle \left|L
ight
angle + \left|\uparrow\right\rangle \left|R
ight
angle + \left|\downarrow\right\rangle \left|R
ight
angle$$

is not entangled because it can be expressed as:

```
(|\uparrow\rangle + |\downarrow\rangle)(|L\rangle + |R\rangle).
```

Example 2 (our entangled state):

$$|\uparrow\rangle |L\rangle + \downarrow |L\rangle |R\rangle.$$

The state of the whole system is better defined than its components:

$$|\uparrow\rangle + |\downarrow\rangle$$
 and $|L\rangle + |R\rangle$.

⁵E. Schrödinger, Naturwissenschaften **23**, 807 (1935)

What can we do with our Schrödinger-cat state?

Spin separation has already been done:

- Left dot \Longrightarrow Spin-up
- Right dot \implies Spin-down



We can do more:

If we rotate spin in the right dot by an angle $\varphi = \pi$ while keeping spin in the left dot unchanged we **arrive at spin oriented upwards in the entire nanodevice**.

Spin rotation in the right dot

Let us see what SOI we can obtain:

$$\hat{H}_{\rm so} = \frac{\alpha_{\rm so}|\boldsymbol{e}|}{\hbar} \left(\boldsymbol{E}_{z} \sigma_{y} - \boldsymbol{E}_{y} \sigma_{z} \right) \hat{\boldsymbol{p}}_{\rm x}.$$
(12)

So far the SOI was generated using lateral gates that create the electric field E_{γ} .



We can also create the electric field E_z by applying a voltage across the substrate and the top gate. In such a case, the SOI is described by

$$\hat{H}_{\rm so} = \frac{\alpha_{\rm so}|\boldsymbol{e}|}{\hbar} \boldsymbol{E}_{\boldsymbol{z}} \sigma_{\boldsymbol{y}} \hat{\boldsymbol{p}}_{\boldsymbol{x}}.$$
 (13)

Spin rotation in the right dot



⁶obrot.exe

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Spin rotation in the right dot

So far, the initial spin was in parallel to the x axis, thus it was an equally weighted linear combination of spin up and down along the z axis. Now, let us assume a **non-equally** weighted linear combination⁷:





⁷inny.exe obrot1.exe

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Complete spin initialization procedure

Initialization can be performed faster if we rotate spins in **both halves** of the nanodevice (not just one) **in antiphase**. The resulting spin will be oriented along the x axis.



- Regardless of the initial spin orientation at the end of the procedure the spin is oriented along the *x* axis.
- The initialization time does not exceed 60 ps, which is significantly less than the spin coherence time (about 20 ns)

Yet another nanodevices



- All employ the Rashba spin-orbit interaction for spin initialization;
- planar nanostructures instead of a quantum wire;
- slightly different procedure due to technological limitations;
- slower initialization time;
- better integration with existing electronics and scalability.

To sum everything up...

Stage 1 (spin separation)







Computational methods

The time evolution of electron's wavefunction was calculated by solving the time-dependent Schrödinger equation

$$i\hbar \frac{\partial}{\partial t} \Psi(x,t) = \hat{H}(x,t)\Psi(x,t)$$

 $\hat{H}(x,t) = -\frac{\hbar^2}{2m^*} \frac{\partial^2}{\partial x^2} + U(x,t) + \hat{H}_{so}(x,t)$

The potential in the nanodevice was calculated using the Poisson equation

$$abla^2 \varphi(\mathbf{r},t) = rac{
ho(\mathbf{r},t)}{arepsilon_0 arepsilon} \implies U(\mathbf{r},t) = -|e|\varphi(\mathbf{r},t)$$

Due to the time-varying voltages applied to the gates, the Poisson equation is solved for every time-step.

Summary

- Simulations were performed for a concrete (described) realistic nanodevice design.
- They take into account all important material parameters.
- The computational methods were tested with nanodevices for which experimental data were available and the results could be compared.
- Good consistency of our results with experimental data allows us to believe that our nanodevices will work according to our predictions.
- Our nanodevice design allows for an ultrafast spin initialization.

Conclusion: It is high time the experimental research in this area was resumed. **Authors' literature:**

- S. Bednarek et al., Ultrafast Spin Initialization in a Gated InSb Nanowire Quantum Dot, Phys. Rev. Applied 11, 034012 (2019)
- S. Bednarek et al., *All-electric single electron spin initialization*, *New J. Phys.* **19** 123006 (2017)
- Pawłowski et al., Phys. Rev. B 96, 115308 (2017)